

A possible enhancement of nuclear fission in scattering with low energy charged particles

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New mechanisms for tunneling enhancement

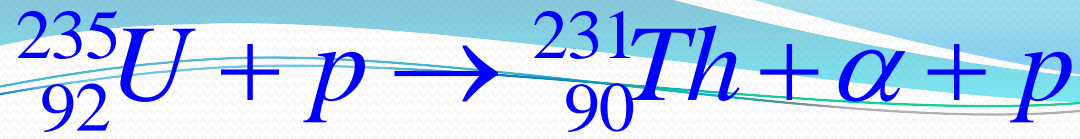
- Tunneling in scattering with multiple degrees of freedom [G. F. Bonini, A. G. Cohen, C. Rebbi and V. A. Rubakov, 1999]
- Alpha-decay induced by collision with low energy protons [B. Ivlev and V. G., 2004]
- Influence of intrinsic structure of composed particles on the tunneling [C. A. Bertulani, V. V. Flambaum and V. G. Zelevinsky, 2007]
-
- Could be important for sub-barrier and close-to-barrier nuclear fission.

Semi-classical approach

$$W \sim \exp(-A(E))$$

$$A(E) = \frac{2}{\hbar} \int dx \sqrt{2m[V(x) - E]}$$

$$A \sim \frac{V}{\hbar\omega} \gg 1$$



$$E = 4.678\text{MeV}$$

Numerical 2-dim: $\Delta E = 0\text{MeV}$ for $\varepsilon_0 \approx 1.85\text{MeV}$ when $\phi \approx 30^\circ$

"Resonance":

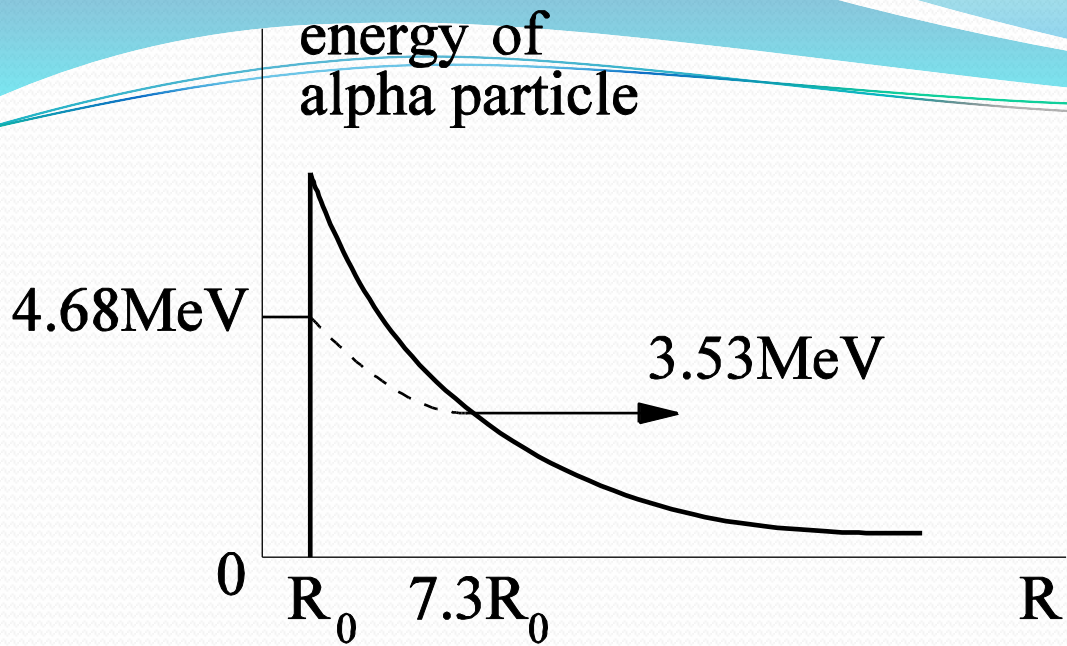
$$\varepsilon_R \approx 0.25\text{MeV} \text{ with } \Delta E = -1.15\text{MeV} \text{ and } \phi \approx 11^\circ.$$

Then

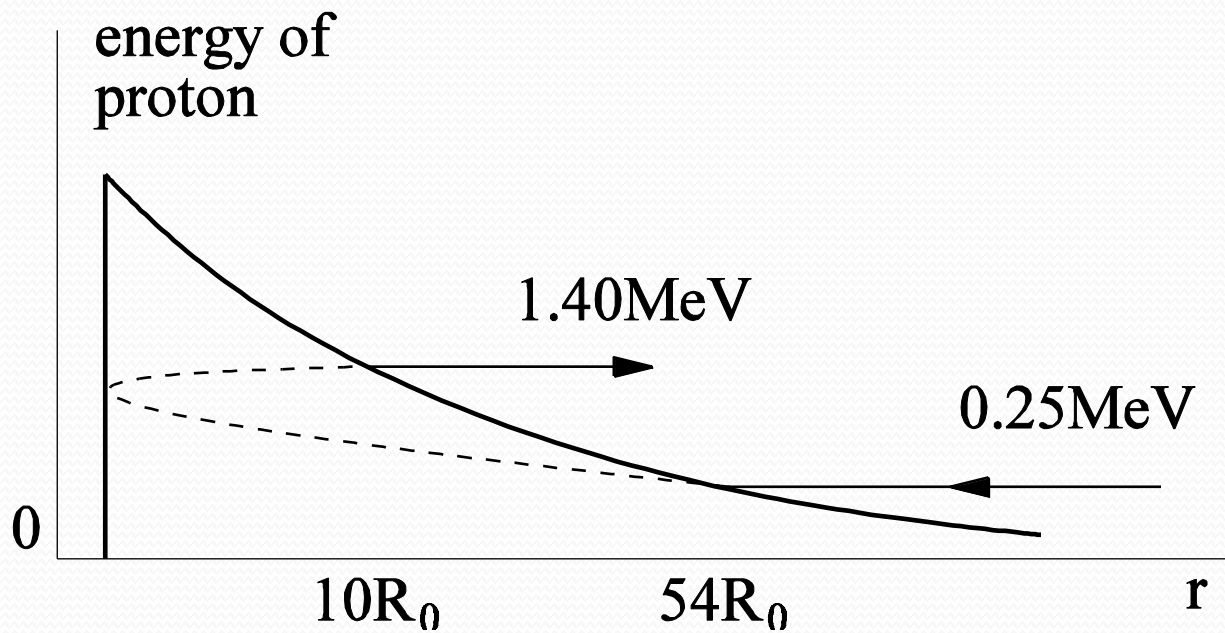
$$E - |\Delta E| \approx 3.53\text{MeV} \quad \text{and} \quad \varepsilon + |\Delta E| \approx 1.40\text{MeV}$$

The shape of the pick

$$\sim \exp\left(-\frac{\varepsilon - \varepsilon_R}{\Delta\varepsilon}\right) \quad \text{with } \Delta\varepsilon \approx 3\text{KeV}$$



(a)



(b)

Fission

$$i\hbar \frac{\partial \Psi(x,t)}{\partial t} = (H_0 + V_t) \Psi(x,t)$$

$$V_t = \frac{eZ}{R(t)} \quad \rightarrow \quad V_\xi \quad \text{with} \quad r(\xi) = a(\sigma \cosh \xi + 1)$$

$$t = \gamma(\sigma \sinh \xi + \xi); \quad \gamma = \sqrt{\frac{m_p a^3}{eZ}}; \quad a = \frac{eZ}{2E_p}$$

$$\sigma = \sqrt{1 + \frac{2E_p \hbar^2 l(l+1)}{m_p (eZ)^2}}$$

Proton potential

$$V_t = \frac{A_1}{|\vec{r}_1 + \vec{R}(t)|} + \frac{A_2}{|\vec{r}_2 + \vec{R}(t)|}; \quad |R| \gg |r_1|, |r_2|$$

\Rightarrow

$$V_t \simeq \frac{(A_1 + A_2)}{R} + \frac{(A_2 r_2 - A_1 r_1)}{R^2} \cos \varphi + \frac{(A_1 + A_2)}{2R^3} (r_1^2 + r_2^2) + \dots$$

Toy model

$$i\hbar \frac{\partial \psi}{\partial t} = \frac{1}{2} \frac{\partial^2 \psi}{\partial x^2} + \frac{1}{2} \omega^2(t) x^2 \psi$$

where

$$\omega^2(t) = \omega_0^2 + \frac{eZ}{4a^3 m_p (\sigma \cosh \xi + 1)^3} = \omega_0^2 (1 + \eta / (\sigma \cosh \xi + 1)^3)$$

$$\eta = \frac{2}{(\alpha Z)^2} \left(\frac{E_p}{\omega_0 \hbar} \right)^2 \left(\frac{E_p}{m_p c^2} \right) \sim 10^{-4} - 10^{-5} \quad \text{for } E_p \sim 1 \text{ MeV}$$

Quantum parametric resonance

$$\omega^2(t) = \omega_0^2(1 + 2\varepsilon \sin((2 + \delta)\omega_0 t)) \quad \text{at} \quad |\varepsilon|, |\delta| \ll 1$$

Fourier Series:

$$t \in [-\pi / \omega_0, \pi / \omega_0] \quad \text{then} \quad \frac{\pi}{\omega_0 \gamma} \geq |\sigma \sinh \xi + \xi|$$

$$\omega_0 \gamma = (\alpha Z) \left(\frac{\omega_0 \hbar}{E_p} \right) \sqrt{\frac{m_p c^2}{E_p}} \sim 10 - 10^2 \gg 1$$

then $\xi \gg 1$

- (L. P. Pitaevsky, A. M. Perelomov, Ya. B. Zeldovich)

QPR-2

$$\varepsilon \sim \frac{1}{(\alpha Z)^2} \left(\frac{E_p}{\omega_0 \hbar} \right)^2 \left(\frac{E_p}{m_p c^2} \right) \sim 2 \cdot (10^{-3} - 10^{-4}) \text{ for } E_p \sim 1 \text{ MeV}$$

$$\Gamma \sim \varepsilon \omega_0 \sim 0.2 \text{ KeV} - 2 \text{ KeV}$$

QPR-barrier penetration

For $|\varepsilon| \sim |\delta| \ll 1$

$$T \sim \frac{\varepsilon^2 \omega_B^2 t^2 / 4}{1 + \varepsilon^2 \omega_B^2 t^2 / 4}$$

$$t \sim D / v_p = (eZ / E_p) / (E_p / 2m_p) \quad \text{and} \quad \omega_B = E_B / \hbar$$

\Rightarrow

$$\omega_B t \sim 2(\alpha Z) \frac{E_B (m_p c^2)}{E_p^2}$$

Then:

$$\varepsilon \omega_B t \sim \frac{2}{(\alpha Z)} \left(\frac{E_p}{E_B} \right) \sim \frac{1}{2} \sim 1 \quad \text{and} \quad T \sim 0.1 - 0.5$$

- (L. P. Pitaevsky, A. M. Perelomov, Ya. B. Zeldovich)

Conclusion

- The possible manifestation of parametric resonance enhancement of fission could be an observation of narrow $(0.2\text{KeV} - 2\text{KeV})$ but very strong resonances induced by collision with low energy protons.